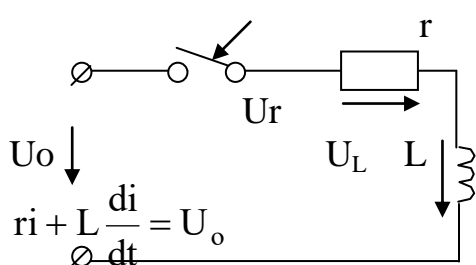


24-ma`ruza. Oddiy elektr zanjirlarida o`tkinchi jarayonlarni klassik usulda xisoblash.

Reja

1. O`tkinchi jarayonni qisoblashning klassik usuli.
2. Kondensatorning aperiodik, chegaraviy aperiodik va tebranma zaryadsizlanishi.

r,l elementlari ketma-ket ulangan zanjirni o`zgarmas kuchlanish manbaiga ulash.



Zanjir $t=0$ da o`zgarmas kuchlanish U_0 ga ulanadi
Zanjirning differensial tenglamasi:

Unga mos bir jinsli tenglama quyidagi ko`rinishga ega:

$$L \frac{di_{\text{apk}}}{dt} + ri_{\text{apk}} = 0$$

i_{apk} aniqlaymiz, uning xarakteristik tenglamasi quyidagicha bo`ladi:

$$L\alpha + r = 0$$

U birgina xaqiqiy va manfiy $\alpha = -\frac{r}{L}$

$$i_{\text{apk}} = Ae^{\alpha t} = Ae^{-\frac{r}{L}t} = Ae^{-\frac{t}{\tau}}$$

bunda: $\alpha = -\frac{r}{L}$ so`nish koeffisienti

Integrallash doimiysini A boshlang`ich shartlar orqali aniqlaymiz. Kommutatsiyaga qadar zanjirdan tok o`tmagan:

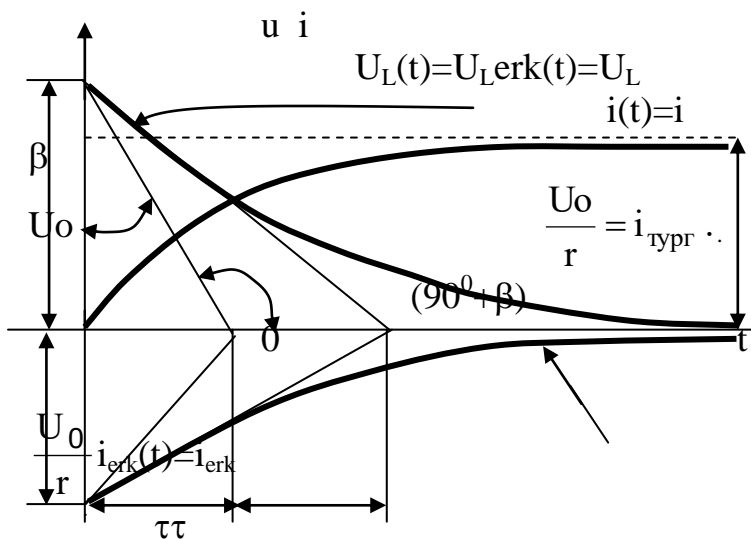
$$i_-(0) = 0$$

demak zanjirni ulash paytida xam u nolga teng bo`lgan $i(0) = 0$ Zanjirda o`tkinchi protsess tugagandan so`ng, faqat r ning qarshiligi bilan aniqlanadigan $i_{\text{typr}} = \frac{U_0}{r}$ tok bo`ladi, unda:

$$i(0) = i_{\text{typr}}(0) + i_{\text{epk}}(0) = \frac{U_0}{r} + A = 0$$

Shunday qilib $A = -\frac{U_0}{r}$ bo`lib o`tkinchi tok i esa:

$$i = i_{\text{typr}} + i_{\text{epk}} = \frac{U_0}{r} - \frac{U_0}{r} e^{-\frac{t}{\tau}} = \frac{U_0}{r} (1 - e^{-\frac{t}{\tau}})$$



τ -vaqt doimiysi; $\tau = L/r$.

Zanjirning ayram elementlaridagi o`tkinchi kuchlanishlarni aniqlash mumkin:

$$U_r = ri = U_0 (1 - e^{-\frac{t}{\tau}})$$

$$U_L = L \frac{di}{dt} = \frac{L}{r} U_0 \left[0 - \left(-\frac{1}{\tau}\right) e^{-\frac{t}{\tau}} \right] = U_0 e^{-\frac{t}{\tau}}$$

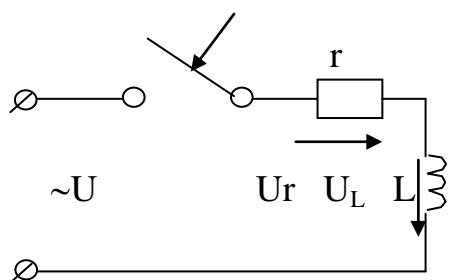
Masala:

Berilgan: $r=50 \text{ Om}$, $L=0,1 \text{ Gn}$, $U_0=200 \text{ B}$.

$$\tau = \frac{L}{r} = 0,002 \text{ cek} ; \quad i_{\text{turg`}} = I_0 = \frac{U_0}{r} = \frac{200}{50} = 4 \text{ A} ; \quad t=5\tau = 0,01 \text{ cek.}$$

$$i = I_0 (1 - e^{-\frac{t}{\tau}}) = 4(1 - e^{-5}) = 4 \cdot 0,099325 = 3,97 \text{ A}$$

R,L elementlari ketma-ket biriktirilgan zanjirni sinusoidal o'zgaruvchan kuchlanish manbaiga ulash.



r va L sinusoidal kuch $U=U_m \sin(\omega t + \psi_u)$

$$L \frac{di}{dt} + ri = U_m \sin(\omega t + \psi_u) \text{ bo'lib}$$

turg'unlashgan tok:

$$i_{\text{typr}} = \frac{U_m}{Z} \sin(\omega t + \psi_u - \varphi) \text{ ni xosil qiladi.}$$

bu yerda: $Z = \sqrt{r^2 + \omega^2 L^2}$ – zanjirning to'la qarshiligi

$\varphi = \arctg \frac{\omega L}{r}$ – kuchlanish u va tok i_{turg} orasidagi siltish burchagi.

Oldingi paragrafdan ma'lumki, tok=erkin tashkil etib:

$$i_{\text{эпк}} = A e^{-\frac{t}{\tau}}$$

Bunda zanjirdagi o'tkinchi tok qo'yidagicha o'zgaradi:

$$i = i_{\text{typr}} + i_{\text{эпк}} = \frac{U_m}{Z} \sin(\omega t + \psi_u - \varphi) + A e^{-\frac{t}{\tau}}$$

A -ni aniqlash uchun boshlang'ich shartlarga murojat qilamiz, kommutatsiya birinchi qonuniga ko'ra:

$$i_-(0) = (0) = \frac{U_m}{Z} \sin(\psi_u - \varphi) + A = 0 \text{ ga egamiz}$$

$$\text{bundan: } A = -\frac{U_m}{Z} \sin(\psi_u - \varphi)$$

$$i = \frac{U_m}{Z} \sin(\omega t + \psi_u - \varphi) - \frac{U_m}{Z} \sin(\psi_u - \varphi) e^{-\frac{t}{\tau}} \text{ bo'ladi.} \quad (1)$$

Induktivlik g'altagining qismalaridagi o'tkinchi kuchlanish qo'yidagicha:

$$U_L = L \frac{di}{dt} = \frac{\omega L}{Z} U_m \cos(\omega t + \psi_u - \varphi + \frac{\pi}{2}) + \frac{r}{Z} U_m \sin(\psi_u - \varphi) e^{-\frac{t}{\tau}} =$$

$$= U_m \left[\sin \varphi \cos(\omega t + \psi_u - \varphi + \frac{\pi}{2}) + \cos \varphi \sin(\psi_u - \varphi) e^{-\frac{t}{\tau}} \right] \quad (2)$$

Induktivlikdagi o`tkinchi tok va kuchlanish uchun olingan ifodalarning to`g`riligiga vakt t ga 0 dan ∞ gacha chegaraviy qiymatlar berib ishonch qosil qilish mumkin. Masalan, vaqt $t=0$ da (1) ning o`ng qismi nolga aylanadi, bu esa induktivlikdagi tok uchun kommutatsiya qonuni tasdiqlaydi:

$$i_-(0) = i(0) = 0$$

O`tkinchi protsess tugagandan so`ng $t=\infty$ da:

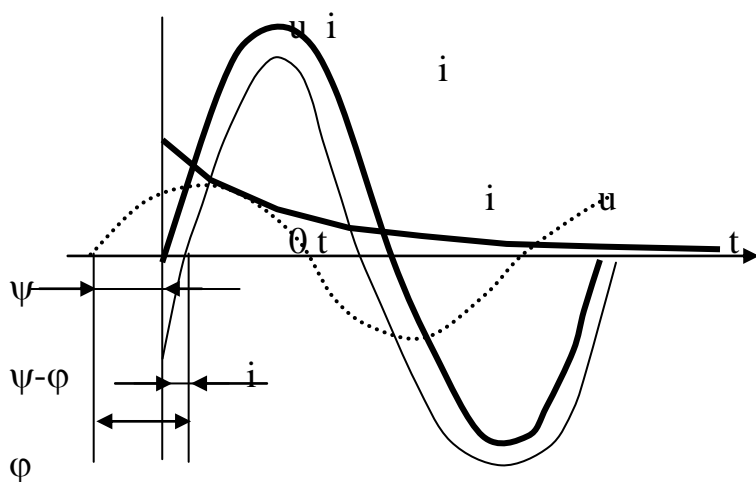
$$i = i_{\text{typ}} = \frac{U_m}{Z} \sin(\omega t + \psi_u - \varphi)$$

chunki $i_{\text{erk}}=0$ 2- tenglamaga ko`ra kommutatsiya paytida:

$$U_L(0) = U_m [\sin \varphi \cos(\psi_u - \varphi) + \cos \varphi \sin(\psi_u - \varphi)] = U_m \sin \psi$$

bo`ladi, ya`ni $t=0$ g`altak qismlardagi kuchlanish zanjirga berilgan kuchlanishning oniy qiymatiga teng.

O`tkinchi tok va o`tkinchi kuchlanish egri chiziq-lari

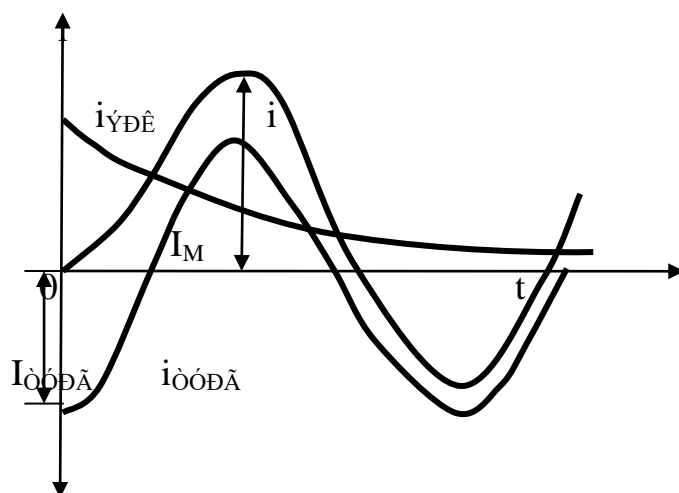


a)

a) rasmda $\psi > 0$, $\varphi > 0$ va $\psi - \varphi < 0$ xollar uchun zanjirga berish kuchlanish $u = U_m \sin(\omega t + \varphi)$ ning (punktir chiziq) va o`tkinchi tok i ning (yo`g`on chiziq) egri

chiziq ko`rsatilgan $i_{\text{opk}} = -\frac{U_m}{Z} \sin(\psi_u - \varphi) e^{-\frac{t}{\tau}}$ mikdori va yo`nalish

kommutatsiya paytida i_{turg} ning boshlang`ich fazasiga bog`liq $\psi_i = \psi_u - \varphi$.



á)

Agar kuchlanishining fazasi ψ_U siltish burchagi φ ga teng bo'lib qolsa, u xolda o'tkinchi tok i faqat turg'un tok i_{TURG} ning qiymatidan iborat bo'ladi, chunki 1-tenglamaga asosan $i_{ERK} = 0$ bo'ladi. Shunda zanjir ulanganda tebranishsiz

sinusoidal tok $i = \frac{U_m}{Z} \sin \omega t$ turg'unlashadi. Agar zanjirni ulash $\psi = \varphi \pm \frac{\pi}{2}$ da sodir bo'lsa, erkin tokning boshlang'ich chayqalishi o'zining maksimum qiymatiga

erishadi $i_{\text{epk}}(0) = \pm \frac{U_m}{Z}$ (b-rasm).

Shu tufayli zanjirga berilgan kuchlanish davrining yarimiga teng bo'lgan vaqtdan so'ng, o'tkinchi tok xam o'zining eng katta qiymati i_{max} -ga erishadi. Ammo nazariy jixatdan tokning bu maksimumi xatto vaqt doimiysi τ ning eng katta qiymatlarida xam turg'unlashgan tok amplitudaviy qiymatining ikkilanganidan ortmaydi.

R.1 elementlari ketma-ket ulangan zanjirda qisqa tutashuv.

Boshlang'ich toki $i(0) \neq 0$ noldan farq qiladigan r , L elementi ketma-ket ulangan zanjir qisqa tutashganda vujudga keluvchi o'tkinchi protsesslarni ko'ramiz. Zanjir kommutatsiyaga qadar qo'shimcha rezistor r_0 orqali o'zgarmas kuchlanish U_0 ga ulangan $t=0$ da qisqa tutashadi deb faraz qilamiz.

Kommutatsiyaga qadar turg'un tok:

$$I_0 = \frac{U_0}{r_0 + r} \text{ bo'lsa,}$$

kommutatsiyadan so'ng nolga teng, chunki kontur manbadan ajratilgan bo'lib, ilgari g'altakka tok I_0 olib kirgan magnit maydonning energiyasi

$$W_m = \frac{LI_0^2}{2}$$

rezistor r da issiqlik energiyasiga aylanib, asta sekin nolgacha kamayadi. O`tkinchi rejimdagi zanjirning differensial tenglamasi:

$$L \frac{di_{\text{эпк}}}{dt} + ri_{\text{эпк}} = 0$$

bo`ladi, o`tkinchi tok esa: $i = i_{\text{эпк}} = Ae^{-\frac{t}{\tau}}$

Integrallash doimiysi Ani $i_{-}(0) = i(0)$ shartidan topamiz, ya`ni $\frac{U_0}{r_0 + r} = A$

$$\text{Shunday qilib : } i = \frac{U_0}{r_0 + r} e^{-\frac{t}{\tau}} = I_0 e^{-\frac{t}{\tau}}$$

Induktiv g`altakdagi o`tkinchi kuchlanish:

$$U_L = L \frac{di}{dt} = -\frac{U_0 r}{r_0 + r} e^{-\frac{t}{\tau}} = -r I_0 e^{-\frac{t}{\tau}}$$

Ya`ni rezistordagi kuchlanish tushuvi:

$$U_r = ri = r I_0 e^{-\frac{t}{\tau}}$$

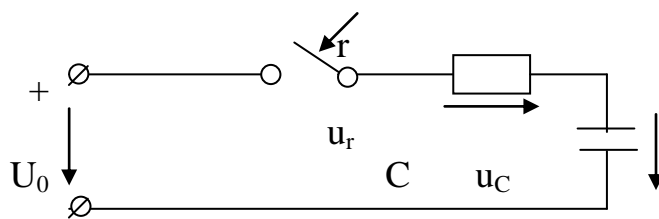
ga teng bo`lib ishorasi qarama-qarshi kommutatsiya paytida o`tkinchi tokning avvalgi qiymati:

$$i_{-}(0) = i(0) = I_0$$

ni saqlagan tok i dan farqli induktiv g`altakdagi kuchlanishi $u_{L-}(0)=0$ dan $U_L(0) = -r I_0$ gacha sakrab o`tishi mumkin (b-chizma).

G`altakdagi e.yu.k. $e_L = -L \frac{di}{dt}$ ning qo`yib ishoralari kuchlanishga va tokning vaqt

bo`yicha o`zgarishga bog`liq. Agar induktiv tarmoqdagi tok orta borsa e.yu.k. ning ishorasi $e_L = -u_L$ manfiy, kamaya borsa musbat bo`ladi. Shunday qilib induktiv g`altakni o`zgarimas kuchlanishni manbai ulash va ajratish tokning qar qanday o`zgarishi e.yu.k. xosil qilishga va uning zanjirdagi boshlangich muvozanat xolatini saqlab turishiga sabab bo`ladi.



Chizmada vaqtning $t=0$ paytida r, C elementlari ketma-ket ulangan zanjirni o'zgaras kuchlanish U_0 manbaiga ulash sxemasi ko'rsatilgan. Sig'imdagi kuchlanish U_C ni o'zgaruvchan xisoblab zanjirning differensial tenglamani tuzamiz:

$$rC \frac{dU_C}{dt} + U_C = U_0$$

bunda: $i = i = \frac{dq}{dt} = C \frac{dU_C}{dt}$ kondensator S qoplamalaridagi o'tkincha tok.

Sig'imdagi o'tkinchi kuchlanish $U_C = U_{\text{turg}} + U_{\text{erk}}$

$$rC \frac{dU_{C_{\text{эпк}}}}{dt} + U_{C_{\text{эпк}}} = 0$$

Bunga mos xarakteristik tenglama:

$$rC\alpha + 1 = 0$$

dan chastota o'lchash $\frac{1}{\text{сек}}$, $\alpha = -\frac{1}{rC}$ ni topamiz.

So'nish koeffisienti α -ning ishorasiga ko'ra,

$$U_{C_{\text{эпк}}} = Ae^{\alpha t} = Ae^{-\frac{t}{rC}}$$

Integrallash doimiysi A ni boshlangich shartlardan topamiz:

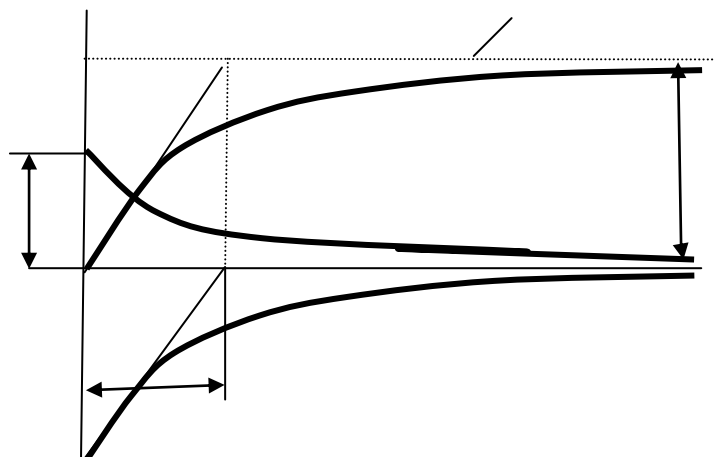
$$\begin{aligned} U_{C_{\text{т}}}(0) &= U_C(0) \\ U_{C_{\text{т}}}(0) = 0 &= U_{C_{\text{т}}}(0) + U_{C_{\text{эпк}}}(0) = U_0 + A \\ &\text{yoki} \\ A &= -U_0 \end{aligned}$$

Kommutatsiyaga qadar sig'imdagi kuchlanish nolga teng $U_C(0)=0$. O'tkinchi protsess tugagandagi turg'unlashgan kuchlanish $U_{\text{Sturg}}=U_0$, chunki o'zgaras tokda kondensatorning qarshiligi ∞ (tok o'tmaydi) zanjirdagi tok $i_{\text{turg}}=0$.

Sig'imdagi kuchlanish:

$$U_C = U_{C_{\text{т}}}(0) + U_{C_{\text{эпк}}} = U_0 - U_0 e^{-\frac{t}{\tau}} = U_0 (1 - e^{-\frac{t}{\tau}})$$

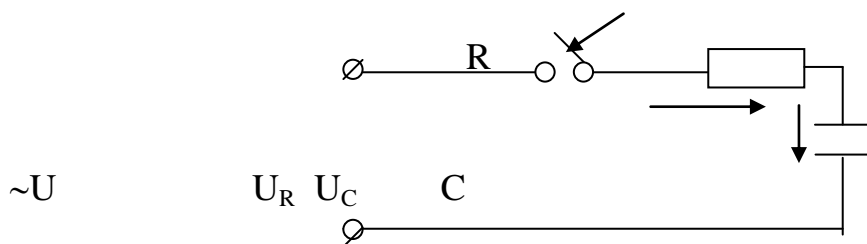
$\tau=rC$ - o'tkinchi protsessning intensivligini xarakterlovchi vaqt doimiysi (sek), u qanchalik katta bo'lsa S ning zaryadlanishi shunchalik sekin boradi va aksincha.



O`tkinchi tok $i = C \frac{dU_C}{dt}$ binoan:

$$i = CU_0 \left[0 - \left(-\frac{1}{rC} e^{-\frac{t}{\tau}} \right) \right] = \frac{U_0}{r} e^{-\frac{t}{\tau}}$$

R,c elementlari ketma-ket ulangan zanjirni sinusoidal o`zgaruvchan kuchlanish manbaiga ulash.



r C elementlar ketma-ket ulangan zanjir $U = U_m \sin(\omega t + \varphi_u)$ manbaga ulanganda o`tkincha kuchlanish va toklar o`zgarish qoqunlarini aniqlaymiz. O`tkinchi kuchlanish:

$$U_C = U_{C\text{typ}} + U_{C\text{эпк}}$$

Bularning birinchisi $i_{\text{тург}}$ - bog`liq ikkinchisi esa $u_{C\text{эпк}} = Ae^{-\frac{t}{\tau}}$ - o`tkinchi.

$$i_{\text{тыпр}} = \frac{U_m}{Z} \sin(\omega t + \psi_u - \varphi) \text{ bo`ladi}$$

áunda: $Z = \sqrt{r^2 + \left(\frac{1}{\omega c}\right)^2}$ - zanjirni to`la qarshiligi,

$\varphi = \arctg\left(-\frac{1}{r\omega c}\right)$ - o`tkinchi protsess tugagandan keyingi u \hat{a} i orasidagi siljish burchagi.

$$U_{C\text{ry pr}} = \frac{U_m}{z\omega C} \sin(\omega t + \psi_u - \varphi - \frac{\pi}{2})$$

chunki u tokdan $\frac{\pi}{2}$ burchakka orqada qoladi, unda:

$$U_C = \frac{U_m}{z\omega C} \sin(\omega t + \psi_u - \varphi - \frac{\pi}{2}) + Ae^{-\frac{t}{\tau}}, \quad \tau = rC$$

Integrallash doimiysi A ni boshlangich shartlardan aniqlanadi:

$$U_C(0) = 0 = \frac{U_m}{z\omega C} \sin(\psi_u - \varphi - \frac{\pi}{2}) + A$$

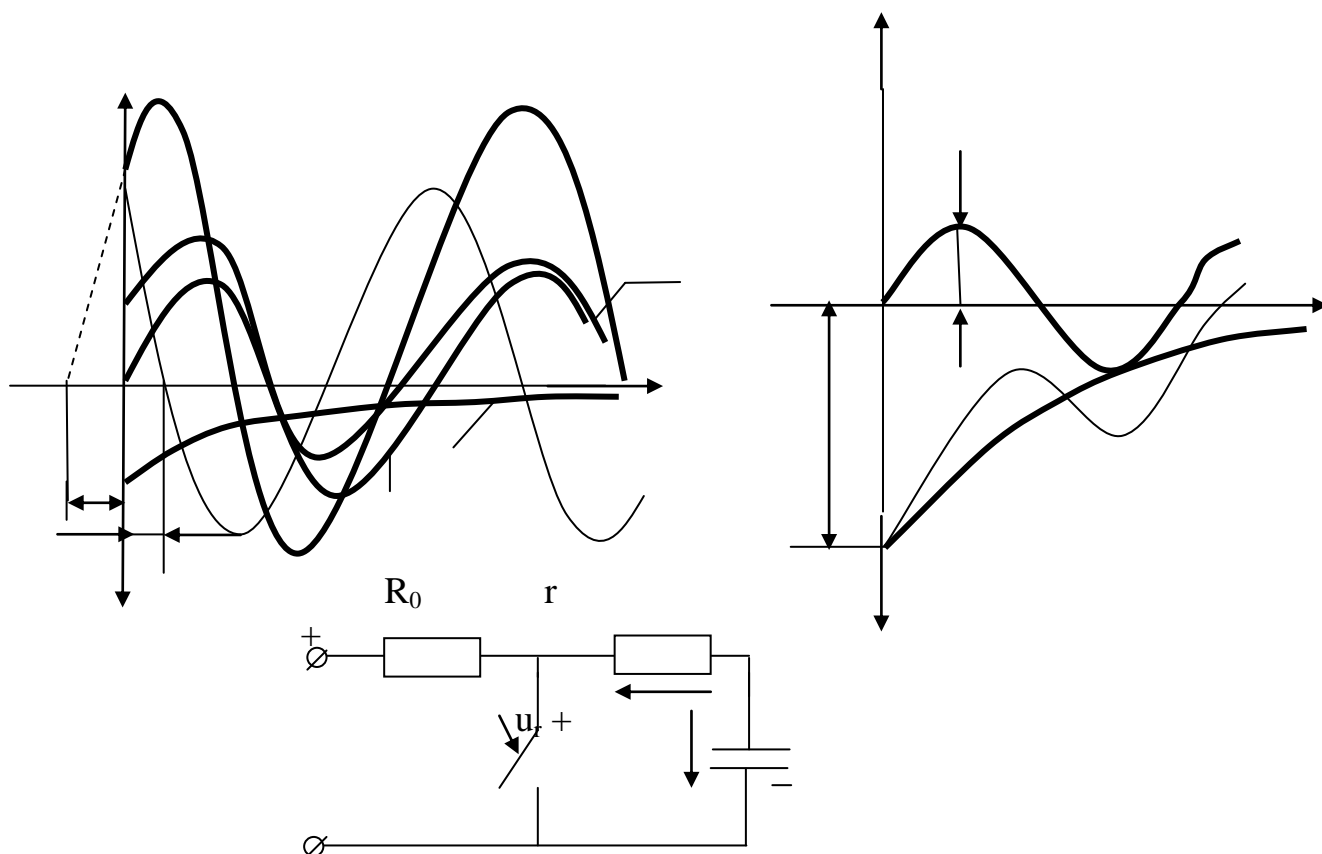
bu yerdan: $A = -\frac{U_m}{z\omega C} \sin(\psi_u - \varphi - \frac{\pi}{2})$

U_S - kuchlanishning o'zgarish qonuni quyidagicha:

$$U_C = \frac{U_m}{z\omega C} \sin(\omega t + \psi_u - \varphi - \frac{\pi}{2}) - \frac{U_m}{z\omega C} \sin(\psi_u - \varphi - \frac{\pi}{2}) e^{-\frac{t}{\tau}}$$

Bundan o'tkinchi tok $i = C \frac{dU_c}{dt}$ ni topamiz:

$$i = \frac{U_m}{z} \sin(\omega t + \psi_u - \varphi) + \frac{U_m}{z} \frac{1}{r\omega C} \sin(\psi_u - \varphi - \frac{\pi}{2}) e^{-\frac{t}{rC}}$$



$$u_0 \quad C$$

$$U_c = U_{c\text{эпк}} = A e^{-\frac{t}{\tau}} \quad \tau = rC$$

Integrallash doimiysi boshlang'ich shartlar yordamida aniqlanadi:

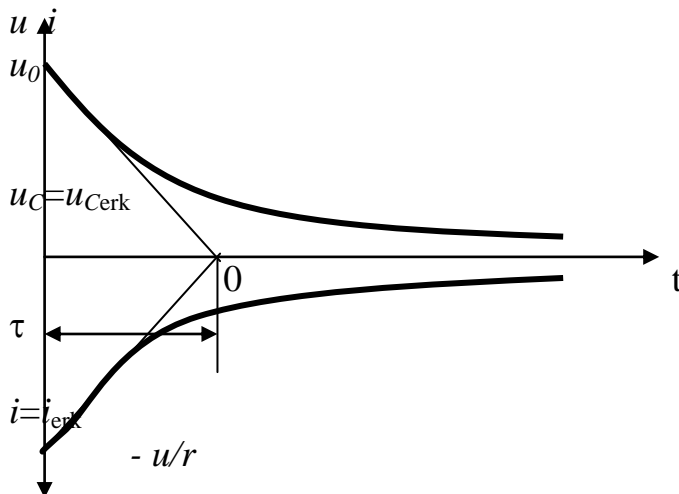
$$u_c(0) = u_c(0) = u_0$$

$$u_c(0) = A = u_0 \text{ bo'ladi yoki}$$

$$u_c = u_0 e^{-\frac{t}{\tau}}$$

Zanjirdagi o'tkinchi tok:

$$i = i_{\text{erk}} = C \frac{du_c}{dt} = -\frac{u_0}{r} e^{-\frac{t}{\tau}}$$

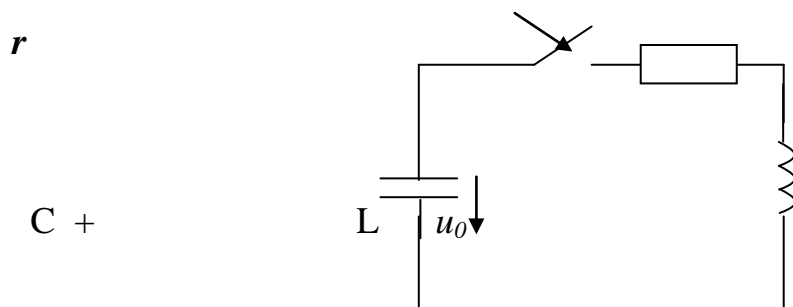


O'tkinchijarayon ($t=0$ dan $t=\infty$ gacha) davomida qarshilik R

$$w_r = \int_0^{\infty} i^2 r dt = -\frac{u_0^2}{r} \int_0^{\infty} e^{-\frac{2t}{\tau}} dt = -\frac{u_0^2}{r} \cdot \frac{rC}{2} \left| e^{-\frac{2t}{\tau}} \right|_0^{\infty} = \frac{Cu_0^2}{2}$$

Kondensatorning zaryadsizlanishi unda to'plangan elektr energiyasi miqdoriga bog'liq

$$W_3 = \frac{Cu_0^2}{2}$$



Shu. o`tkinchi jarayon tugagandan keyin tok va kuchlanishlarning turg`unlashgan qiymatlari nolga teng. Zanjirning differensial tenglamasi:

$$\frac{d^2 i}{dt^2} + \frac{r}{L} \frac{di}{dt} + \frac{i}{LC} = 0 \quad (1)$$

$$\frac{r}{L} = 2\delta; \quad \omega_0^2 = \frac{1}{LC}$$

$$\alpha^2 + 2\delta\alpha + \omega_0^2 = 0 \quad (2)$$

Uning ildizlari:

$$\alpha_1 = -\delta + \sqrt{\delta^2 - \omega_0^2}$$

$$\alpha_2 = -\delta - \sqrt{\delta^2 - \omega_0^2}$$

$$i = A_1 e^{\alpha_1 t} + A_2 e^{\alpha_2 t}$$

Zanjir o`tkinchi jarayon r, L va C parametrlarning miqdorlari nisbati bilan aniqlanadi. Unda nisbat uch variantda bo`lishi mumkin:

1) α_1 va α_2 ildizlar qariqiy va turlicha, yani $\delta^2 - \omega_0^2 > 0$;

2) α_1 va α_2 ildizlar qariqiy va teng, ya'ni $\delta^2 - \omega_0^2 = 0$;

3) $\delta^2 - \omega_0^2 < 0$ bo`lganda $\alpha_1 = -\delta + j\sqrt{\omega_0^2 - \delta^2}$ } ildizlar qo`shma komplekslardan
 $\alpha_2 = -\delta - j\sqrt{\omega_0^2 - \delta^2}$ iborat.

Dastlabki ikki variantda kondensatorning zaryadsizlanishi aperiodik xarakterda bo`lib, uchinchi variantda esa bu jarayon davriy yoki tebranma bo`ladi.

Sinov savollari.

1. Qisqa tutashuv qolatida g`altakdagi magnitaviy maydon energiyasi qanday energiya turiga aylanadi?
2. O`tkinchi jarayon intensivligini qanday kattalik xarakterlaydi?

3. Sig`imdagi turg`unlashgan kuchlanishni yozib bering?
4. So`nish koeffisienti α qanday aniqlanadi?